The pressure, densities and first order phase transitions associated with multidimensional SOFT

Shmuel Friedland Univ. Illinois at Chicago

Dynamics Seminar, College Park, October 28, 2010

Outline of the talk

- Motivation: Ising model
- Subshifts of Finite Type
- \bigcirc Pressure P_{Γ}
- Density points and density entropy
- Convex functions
- First order phase transition
- The maximum principle
- d-Dimensional Monomer-Dimers
- Friendly colorings
- Computation of pressure



Motivation: Ising model - 1925

On lattice \mathbb{Z}^d two kinds of particles: spin up 1 and spin down 2. Each neighboring particles located on $(\mathbf{i}, \mathbf{i} + \mathbf{e}_j)$ interact with energy -J if both locations are occupied by the same particles, and with energy J if the two sites are occupied by two different particles. In addition each particle has a magnetization due to the external magnetic field. The energy of the particle of type 1 is H while the energy of the particle of type 2 is -H. The energy of $E(\phi)$ of a given finite configuration of particles in \mathbb{Z}^d is the sum of the energies of the above type.

Ferromagnetism J > 0: all spins are up or down.

Antiferromagnetism J < 0 half spins up and down

(Lowest free energy)

Phase transition:

from one state to another as the temperature varies

Energy: $\frac{k}{T}E(\phi)$

Subshifts of Finite Type-SOFT

$$< n > := \{1, 2, 3, ..., n\}$$

ALPHABET ON *n* LETTERS - COLORS.



Coloring of \mathbb{Z}^d in n coloring = Full \mathbb{Z}^d shift on n symbols Example of SOFT: (0-1) LIMITED CHANNEL HARD CORE LATTICE or NEAR NEIGHBOR EXCLUSION $n=2, <2>=\{1,2\}=\{1,0\}$ $(2\equiv0)$. NO TWO 1's ARE NEIGHBORS.



One dimensional SOFT

 $\Gamma \subseteq < n > \times < n >$ directed graph on n vertices $C_{\Gamma}(< m >)$ -all Γ allowable configurations of length m:

$$\{a = a_1 ... a_m = (a_i)_1^m : \langle m \rangle \rightarrow \langle n \rangle (a_i, a_{i+1}) \in \Gamma\}$$

 $C_{\Gamma}(\mathbb{Z})$ -all Γ allowable configurations (tilings) on \mathbb{Z} :

$$\{a = (a_i)_{i \in \mathbb{Z}} : \mathbb{Z} \rightarrow < n >, (a_i, a_{i+1}) \in \Gamma\}$$

Hard core model:

$$n=2, \Gamma=\{\bullet\bullet,\bullet\bullet,\bullet\bullet\}$$





MD SOFT=Potts Models

Dimension $d \geq 2$. For $\mathbf{m} \in \mathbb{N}^d$

$$<$$
 m $> := < m_1 > \times ... \times < m_d >$

$$\operatorname{vol}(\mathbf{m}) := |m_1| \times \ldots \times |m_d|$$

$$\Gamma := (\Gamma_1, \ldots, \Gamma_d), \Gamma_i \subset \langle n \rangle \times \langle n \rangle$$

 $C_{\Gamma}(<\mathbf{m}>)$ -all Γ allowable configurations of \mathbf{m} :

$$a = (a_i)_{i \in \langle \mathbf{m} \rangle} : \langle \mathbf{m} \rangle \rightarrow \langle n \rangle$$

s.t.
$$(a_i, a_{i+e_i}) \in \Gamma_i$$
 if $i, i + e_i \in \langle m \rangle$

$$\mathbf{e}_j = (\delta_{1j}, \ldots, \delta_{dj}), j = 1, \ldots, d.$$

Example:





Γ

12

```
For \phi \in C_{\Gamma}(\langle \mathbf{m} \rangle) - \mathbf{c}(\phi) := (c_1(\phi), \dots, c_n(\phi))
denotes coloring distribution of configuration \phi
c_i(\phi)-the number of times the particle i appears in \phi
\frac{1}{\text{vol(m)}}\mathbf{c}(\phi) \in \Pi_n - coloring frequency of \phi
\Pi_n(\text{vol}(\mathbf{m})) all \mathbf{c} \in \mathbb{Z}_+^n s.t. \frac{1}{\text{vol}(\mathbf{m})}\mathbf{c} \in \Pi_n
C_{\Gamma}(\langle \mathbf{m} \rangle, \mathbf{c}) denotes all \phi \in C_{\Gamma}(\langle \mathbf{m} \rangle) with \mathbf{c}(\phi) = \mathbf{c}.
C_{\Gamma,per}(\langle \mathbf{m} \rangle) \subseteq C_{\Gamma}(\langle \mathbf{m} \rangle)-m-periodic configurations
C_{\Gamma}(\mathbb{Z}^d)-are-\Gamma allowable configurations of \mathbb{Z}^d
Assumption: C_{\Gamma}(\mathbb{Z}^d) \neq \emptyset
u_i \in \mathbb{R} energy of particle i \in \langle n \rangle
\mathbf{u} := (u_1, \dots, u_n) \in \mathbb{R}^n energy vector
E(\phi) = \mathbf{c}(\phi) \cdot \mathbf{u} Energy of configuration \phi
Near neighbor interaction model, can be fit to the above noninteraction
model by considering the coloring of the cube \langle (3,\ldots,3) \rangle as one
particle
```

Similarly short range interaction model

Pressure

Grand partition function

$$\begin{split} &Z_{\Gamma}(\textbf{m}, \textbf{u}) := \sum_{\phi \in C_{\Gamma}(\langle \textbf{m} \rangle)} e^{\textbf{c}(\phi) \cdot \textbf{u}} \\ &\log Z_{\Gamma}(\textbf{m}, \textbf{u}) \text{ subadditive in each component of } \textbf{m} \text{ and convex in } \textbf{u} \\ &\frac{1}{\text{vol}(\textbf{m})} \log Z_{\Gamma}(\textbf{m}, \textbf{u}) - \text{average energy or pressure} \\ &P_{\Gamma}(\textbf{u}) := \lim_{\textbf{m} \to \infty} \frac{1}{\text{vol}(\textbf{m})} \log Z_{\Gamma}(\textbf{m}, \textbf{u}) \\ &\text{Pressure of } \Gamma\text{-SOFT}, \text{ (Pressure of the Potts model)} \\ &h_{\Gamma} := P_{\Gamma}(\textbf{0})\text{-ENTROPY of } \Gamma\text{-SOFT} \\ &P_{\Gamma}(\textbf{u}) \text{ is a convex Lipschitz function on } \mathbb{R}^n \\ &|P_{\Gamma}(\textbf{u}) - P_{\Gamma}(\textbf{v})| \leq ||\textbf{u} - \textbf{v}||_{\infty} := \max |u_i - v_i| \\ &P_{\Gamma}(\textbf{u} + t\textbf{1}) = P_{\Gamma}(\textbf{u}) + t \\ &P_{\Gamma} \text{ has the following properties:} \\ &\text{Has subdifferential } \partial P_{\Gamma}(\textbf{u}) \text{ for each } \textbf{u} \\ &\partial P_{\Gamma}(\textbf{u}) \subseteq \Pi_n \text{ for each } \textbf{u} \\ &\text{Has differentiable } \nabla P_{\Gamma}(\textbf{u}) \text{ a.e.} \end{split}$$

Density points and density entropy

 $\mathbf{p} \in \Pi_n$ density point of $C_{\Gamma}(\mathbb{Z}^d)$ when there exist sequences of boxes $\langle \mathbf{m}_q \rangle \subseteq \mathbb{N}^d$ and color distribution vectors $\mathbf{c}_q \in \Pi_n(\mathrm{vol}(\mathbf{m}_q))$ $\mathbf{m}_q \to \infty, \ C_{\Gamma}(\langle \mathbf{m}_q \rangle, \mathbf{c}_q) \neq \emptyset \ \forall q \in \mathbb{N}, \ \text{and} \ \lim_{q \to \infty} \frac{\mathbf{c}_q}{\mathrm{vol}(\mathbf{m}_q)} = \mathbf{p}$ Π_{Γ} the set of all density points of $C_{\Gamma}(\mathbb{Z}^d)$ Π_{Γ} is a closed set For $\mathbf{p} \in \Pi_{\Gamma}$ the color density entropy $h_{\Gamma}(\mathbf{p}) := \sup_{\mathbf{m}_q, \mathbf{c}_q} \lim\sup_{q \to \infty} \frac{\log \#C_{\Gamma}(\langle \mathbf{m}_q \rangle, \mathbf{c}_q)}{\mathrm{vol}(\mathbf{m}_q)} \geq 0$ where the supremum is taken over all sequences satisfying the above conditions

 h_{Γ} is upper semi-continuous on Π_{Γ}

Convex functions

```
f: \mathbb{R}^n \to [-\infty, \infty] convex.
\operatorname{dom} f := \{ \mathbf{x} \in \mathbb{R}^m : f(\mathbf{x}) < \infty \}
f proper if f: \mathbb{R}^n \to \mathbb{R} := (-\infty, \infty] and f \not\equiv \infty
f closed if f is lower semi-continuous.
q subgradient: f(\mathbf{x}) > f(\mathbf{u}) + \mathbf{q}^{\top}(\mathbf{x} - \mathbf{u}) \ \forall \mathbf{x}
\partial f(\mathbf{u}) \subset \mathbb{R}^n the subset of subgradients of f at \mathbf{u}
ASSUMPTION: f is proper and closed
\partial f(\mathbf{u}) is a closed nonempty set for each \mathbf{u} \in \mathrm{ri} \, \mathrm{dom} \, f
f is differentiable at \mathbf{u} \iff \partial f(\mathbf{u}) = \{\nabla f(\mathbf{u})\}\
diff f - the set of differentiability points of f
\nabla f continuous on diff f and \overline{\text{diff } f} \supset \text{dom } f
The conjugate, (Legendre transform) f^* defined:
f^*(\mathbf{y}) := \sup_{\mathbf{y} \in \mathbb{R}^n} \mathbf{x}^{\top} \mathbf{y} - f(\mathbf{x}) \text{ for each } \mathbf{y} \in \mathbb{R}^m
f^* is a proper closed function and f^{**} = f
```

P_{Γ}^* and color density entropy

```
Thm 1: h_{\Gamma}(\mathbf{p}) \leqslant -P_{\Gamma}^*(\mathbf{p}) \ \forall \mathbf{p} \in \Pi_{\Gamma}.
P_{\Gamma}(\mathbf{u}) = \max_{\mathbf{p} \in \Pi_{\Gamma}} (\mathbf{p}^{\top} \mathbf{u} + h_{\Gamma}(\mathbf{p})), \mathbf{u} \in \mathbb{R}^{n}
\Pi_{\Gamma}(\mathbf{u}) := \operatorname{arg\,max}_{\mathbf{p} \in \Pi_{\Gamma}}(\mathbf{p}^{\top}\mathbf{u} + h_{\Gamma}(\mathbf{p})) = \{\mathbf{p} \in \Pi_{\Gamma} : P_{\Gamma}(\mathbf{u}) = \mathbf{p}^{\top}\mathbf{u} + h_{\Gamma}(\mathbf{p})\}
For each \mathbf{p} \in \Pi_{\Gamma}(\mathbf{u}), h_{\Gamma}(\mathbf{p}) = -P_{\Gamma}^{*}(\mathbf{p}).
\Pi_{\Gamma}(\mathbf{u}) \subseteq \partial P_{\Gamma}(\mathbf{u}).
\mathbf{u} \in \operatorname{diff} P_{\Gamma} \Rightarrow \Pi_{\Gamma}(\mathbf{u}) = \{\nabla P_{\Gamma}(\mathbf{u})\}.
Therefore \partial P_{\Gamma}(\operatorname{diff} P_{\Gamma}) \subseteq \Pi_{\Gamma}.
S(\mathbf{u}), \mathbf{u} \in \mathbb{R}^n \setminus \operatorname{diff} P_{\Gamma}
are all the limits of sequences
\nabla P_{\Gamma}(\mathbf{u}_i), \mathbf{u}_i \in \operatorname{diff} P_{\Gamma} \text{ and } \mathbf{u}_i \to \mathbf{u}.
Then S(\mathbf{u}) \subseteq \Pi_{\Gamma}(\mathbf{u}).
conv \Pi_{\Gamma}(\mathbf{u}) = \text{conv } S(\mathbf{u}) = \partial P_{\Gamma}(\mathbf{u}).
\partial P_{\Gamma}(\mathbb{R}^n) \subseteq \text{conv } \Pi_{\Gamma} \subseteq \Pi_n.
```

conv $\Pi_{\Gamma} = \text{dom } P_{\Gamma}^*$.

Outline of proof

```
From the definitions of P_{\Gamma}(\mathbf{u}), \mathbf{p},
h_{\Gamma}(\mathbf{p}) := \sup_{\mathbf{m}_q, \mathbf{c}_q} \limsup_{q 	o \infty} rac{\log \# C_{\Gamma}(\langle \mathbf{m}_q 
angle, \mathbf{c}_q)}{\operatorname{vol}(\mathbf{m}_q)} \geq 0
P_{\Gamma}(\mathbf{u}) \geqslant \mathbf{p}^{\top}\mathbf{u} + h_{\Gamma}(\mathbf{p}) \Rightarrow P_{\Gamma}(\mathbf{u}) \geqslant \sup_{\mathbf{p} \in \Pi_{\Gamma}} \mathbf{p}^{\top}\mathbf{u} + h_{\Gamma}(\mathbf{p}) \Rightarrow
-h_{\Gamma}(\mathbf{p})\geqslant P_{\Gamma}^{*}(\mathbf{p})\Rightarrow \Pi_{\Gamma}\subseteq \operatorname{dom} P_{\Gamma}^{*}
\mathbf{C}(\mathbf{m}, \mathbf{u}) := \arg \max_{\mathbf{c} \in \Pi_n(\text{vol}(\mathbf{m}))} \# C_{\Gamma}(\langle \mathbf{m} \rangle, \mathbf{c}) e^{\mathbf{c}^{\top} \mathbf{u}}
\textit{Z}_{\Gamma}(\textbf{m},\textbf{u}) = \textit{O}(\text{vol}(\textbf{m})^{\textit{n}-1}) \# \textit{C}_{\Gamma}(\langle \textbf{m} \rangle,\textbf{c}(\textbf{m},\textbf{u})) e^{\textbf{c}(\textbf{m},\textbf{u})^{\top}\textbf{u}}
Let \mathbf{m}_q \to \infty s.t. \frac{\mathbf{c}(\mathbf{m}_q,\mathbf{u})}{\mathbf{vol}(\mathbf{m}_q)} \to \mathbf{p}(\mathbf{u}) \Rightarrow
P_{\Gamma}(\mathbf{u}) \leqslant \mathbf{p}(\mathbf{u})^{\top} \mathbf{u} + \limsup_{q \to \infty} \frac{\log \#C_{\Gamma}(\langle \mathbf{m}_q \rangle, \mathbf{c}(\mathbf{m}_q, \mathbf{u}))}{\operatorname{vol}(\mathbf{m}_r)} \leqslant \mathbf{p}(\mathbf{u})^{\top} \mathbf{u} + h_{\Gamma}(\mathbf{p}(\mathbf{u}))
For p \in \Pi_{\Gamma}(\mathbf{u}) use maximal characterization
P_{\Gamma}(\mathbf{u} + \mathbf{v}) \geqslant \mathbf{p}^{\top}(\mathbf{u} + \mathbf{v}) + h_{\Gamma}(\mathbf{p}) = \mathbf{p}^{\top}\mathbf{v} + P_{\Gamma}(\mathbf{u})
So p \in \partial P_{\Gamma}(\mathbf{u}) \Rightarrow \Pi_{\Gamma}(\mathbf{u}) \subseteq \partial P_{\Gamma}(\mathbf{u}) \Rightarrow
\mathbf{u} \in \operatorname{diff} P_{\Gamma} \Rightarrow \Pi_{\Gamma}(\mathbf{u}) = \{\nabla P_{\Gamma}(\mathbf{u})\}\
```

First order phase transition

Claim: For $\mathbf{u} \in \mathbb{R}^n$ each $\mathbf{p} \in \Pi_{\Gamma}(\mathbf{u})$ is the set of possible density of n colors in an allowable configurations from $C_{\Gamma}(\mathbb{Z}^d)$ with the potential \mathbf{u} . For $\mathbf{u} \in \operatorname{diff} P_{\Gamma} \mathbf{p}(\mathbf{u}) = \nabla P_{\Gamma}(\mathbf{u})$ is a unique density.

Claim: Any point of nondifferentiabity of P_{Γ} is a point of the phase transition.

Proof Let $\mathbf{u} \in \mathbb{R}^n \setminus \operatorname{diff} P_\Gamma$ Then ∂P_Γ consists of more than one point. Thm 1 yields that $\partial P_\Gamma(\mathbf{u}) = \operatorname{conv} S(\mathbf{u}) \subseteq \Pi_\Gamma(\mathbf{u})$. $S(\mathbf{u})$ consists of more than one point. Hence $\Pi_\Gamma(\mathbf{u})$ consists of more than one density for \mathbf{u} . $\mathbf{u} \in \mathbb{R}^n \setminus \operatorname{diff} P_\Gamma$ is called a point of phase transition, or a phase transition point of the first order.

Ergodic Notions

```
C_{\Gamma}(\mathbb{Z}^d)-a compact metric space.
It is invariant under the shifts
\sigma_i: C_{\Gamma}(\mathbb{Z}^d) \to C_{\Gamma}(\mathbb{Z}^d), i = 1, \ldots, d
\sigma_i(\phi) is obtained by shifting the allowable configuration \phi \in C_{\Gamma}(\mathbb{Z}^d)
using the transformation \mathbf{x} \mapsto \mathbf{x} - \mathbf{e}_i.
Let \mathcal{M}_{\Gamma} be the compact set of invariant measures on C_{\Gamma}(\mathbb{Z}^d) with
respect to \sigma_i, i = 1, \ldots, d.
h_{\Gamma}(\mu) -Kolmogorov-Sinai entropy for \mu \in \mathcal{M}_{\Gamma}
h_{\Gamma}(\mu) = \lim_{m \to \infty} \frac{1}{(2m+1)^d} H_{\mu}(\vee_{-m \le i_1, \dots, i_d \le m} \sigma_1^{i_1} \dots \sigma_d^{i_d} \mathcal{A})
where \mathcal{A} = \{A_1, \dots, A_n\}
a cylinder partition of \mathbb{C}_{\Gamma}(\mathbb{Z}^d).
A_i - the set of all configurations \phi \in C_{\Gamma}(\mathbb{Z}^d) s.t. \mathbf{0} \in \mathbb{Z}^d colored by color
i in \phi.
```

The maximum principle

$$f_{\mathbf{u}}: C_{\Gamma} \to \mathbb{R}$$
 be given by $f_{\mathbf{u}}(\phi) = u_i$ for $\phi \in A_i, \mathbf{u} = (u_1, \dots, u_n)$. $P_{\Gamma}(\mathbf{u}) = \max_{\mu \in \mathcal{M}_{\Gamma}} h_{\Gamma}(\mu) + \int f_{\mathbf{u}}(\mathbf{x}) d\mu(\mathbf{x})$ $\mu_{\mathbf{u}} \in \mathcal{M}_{\Gamma}$ is maximal if $P_{\Gamma}(\mathbf{u}) = h_{\Gamma}(\mu) + \int f_{\mathbf{u}}(\mathbf{x}) d\mu(\mathbf{x})$ \mathbf{u} -ergodic phase transition if there are at least two maximal $\mu_{\mathbf{u}}$ measures Conjecture If $\mathbf{u} \in \mathbb{R}^n \setminus \text{diff } P_{\Gamma}$ then \mathbf{u} is an ergodic phase transition Special case studied case in the literature $\mathbf{u} = 0$: The entropy $h_{\Gamma} = P_{\Gamma}(\mathbf{0}) = \max_{\mathbf{p} \in \Pi_{\Gamma}} h_{\Gamma}(\mathbf{p}) = \max_{u \in \mathcal{M}_{\Gamma}} h_{\Gamma}(\mu)$

d-Dimensional Monomer-Dimers

Dimer: $(\mathbf{i}, \mathbf{j}), \ \mathbf{j} = \mathbf{i} + \mathbf{e}_k \in \mathbb{Z}^d$.

any partition of \mathbb{Z}^d to dimers (1-factor).

Monomer: occupies $\mathbf{i} \in \mathbb{Z}^d$.

any partition of \mathbb{Z}^d to monomer-dimers

is 1-factor of a subset of \mathbb{Z}^d .

Dimer and Monomer-Dimer are SOFT

$$0 = \tilde{h}_1 \leq \tilde{h}_2 \leq ... \leq \tilde{h}_d \leq ... (dimers)$$

$$\log \frac{1+\sqrt{5}}{2} = h_1 \le h_2 \le \dots \le h_d \le \dots$$
(monomer – dimer)

Fisher, Kasteleyn and Tempreley 61

$$\tilde{h}_2 = \frac{1}{\pi} \sum_{i=0}^{\infty} \frac{(-1)^i}{(2i+1)^2} = 0.29156090...$$



Hammersley's results

```
Hammersley in 60's studied extensively the monomer-dimer model. He
showed \Pi_{\Gamma} = \Pi_{d+1} for d-dimensional model \mathbf{p} = (p_1, \dots, p_d, p_{d+1})
p_i-the dimer density in \mathbf{e}_i-direction i = 1, \dots, d p_{d+1}-the monomer
density Hammersley studied p := p_1 + ... + p_d-the total dimer density
h_d(p)-the p-dimer density in \mathbb{Z}^d, p \in [0,1]
He showed h_d(p)-concave continuous function on [0, 1]
Heilman and Lieb 72: h_d(p) analytic on (0,1)
No phase transition in parameter p \in (0,1)
Au-Yang and Perk 84: Phase transition at p = 1
Friedland-Krop-Lundow-Markstrom 08
h_d(p) + \frac{1}{2}(p \log p + (1-p) \log(1-p)) concave
```

The Graphs for $h_2(p)$

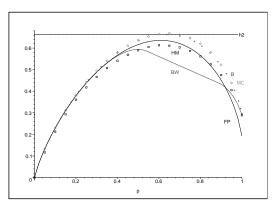
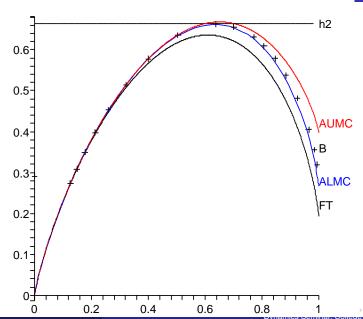


Figure: HM is the lower bound of Hammersley-Menon, BW is the lower bound of Bondy-Welsh, FP is the lower bound of Friedland-Peled, MC is the Monte Carlo estimate of Hammersley-Menon, B are Baxter's estimates, and h2 is the true value of $h_2 = \max h_2(p)$.

Graph



Friendly colorings

Thm 1 implies:

For any Potts model $h_{\Gamma}(\cdot): \Pi_{\Gamma} \to \mathbb{R}_+$ is concave on every convex subset of $\Pi_{\Gamma}(\mathbb{R}^n)$.

To get the exact analog of Hammersley's result

$$\Gamma = (\Gamma_1, \dots, \Gamma_d) \text{ on } \langle n \rangle$$

 $\mathcal{F} = \cup_{\mathbf{m} \in \mathbb{N}^d} C_{\Gamma}(\langle \mathbf{m} \rangle)$, where $C_{\Gamma}(\langle \mathbf{m} \rangle) \subseteq C_{\Gamma}(\langle \mathbf{m} \rangle)$ for each $\mathbf{m} \in \mathbb{N}^d$, friendly: if whenever a box $\langle \mathbf{m} \rangle$ is cut in two and each part is colored by a coloring in \mathcal{F} , the combined coloring is in \mathcal{F} .

Γ friendly if there exist a friendly set $\mathcal{F} = \bigcup_{\mathbf{m} \in \mathbb{N}^d} \widetilde{C}_{\Gamma}(\langle \mathbf{m} \rangle)$ and a constant vector $\mathbf{b} \in \mathbb{N}^d$ such that if any box $\langle \mathbf{m} \rangle$ is padded with an envelope of width b_i in the direction of \mathbf{e}_i , then each Γ-allowed coloring of $\langle \mathbf{m} \rangle$ can be extended in the padded part to a coloring in \mathcal{F} .

Examples of friendly colorings

 Γ has a friendly color $f \in \langle n \rangle$, i.e., for each $i \in \langle d \rangle$ $(f,j), (j,f) \in \Gamma_i$ for all $j \in \langle n \rangle$

Then $C_{\Gamma}(\mathbf{m})$ are Γ -allowed colorings of $\langle \mathbf{m} \rangle$ whose boundary points are colored with f

Hard-core model: $\Gamma_i = \{(1,1), (1,2), (2,1)\}$, has friendly color f = 1.

 $C_{\Gamma}(\langle \mathbf{m} \rangle)$ the set of tilings of $\langle \mathbf{m} \rangle$ by monomers and dimers, i.e., the coverings in which no dimer protrudes out of $\langle \mathbf{m} \rangle$, as in Hammersley

P_{Γ}^* for friendly colorings

Thm 2: Let $\Gamma = (\Gamma_1, \dots, \Gamma_d)$ be a friendly coloring digraph. Then

 Π_{Γ} is convex. Hence $\Pi_{\Gamma} = \operatorname{dom} P_{\Gamma}^*$.

 $h_{\Gamma}(\cdot):\Pi_{\Gamma}\to\mathbb{R}_{+}$ is concave.

For each $\mathbf{u} \in \mathbb{R}^n$, $\Pi_{\Gamma}(\mathbf{u}) = \partial P_{\Gamma}(\mathbf{u})$.

For each $\mathbf{u} \in \mathbb{R}^n$, $h_{\Gamma}(\cdot)$ is an affine function on $\partial P_{\Gamma}(\mathbf{u})$.

 $h_{\Gamma}(\mathbf{p}) = -P_{\Gamma}^{*}(\mathbf{p})$ for each $\mathbf{p} \in \Pi_{\Gamma}$.

Outline of proof

```
(1). Let \alpha \in \widetilde{C}_{\Gamma}(\langle \mathbf{m} \rangle), \mathbf{c}(\alpha) = (c_1, \ldots, c_n) \in \Pi_n(\text{vol}(\mathbf{m})) color frequency vector of \alpha, and \mathbf{p} := \frac{1}{\text{vol}(\mathbf{m})} \mathbf{c}(\alpha). For \mathbf{j} = (k_1, \ldots, k_d) \in \mathbb{N}^d let \mathbf{j} \cdot \mathbf{m} := (k_1 m_1, \ldots, k_d m_d). View \langle \mathbf{j} \cdot \mathbf{m} \rangle as a box composed of vol(\mathbf{j}) boxes isomorphic to \langle \mathbf{m} \rangle color each box by \alpha obtaining a coloring \alpha(\mathbf{j} \cdot \mathbf{m}) \in \widetilde{C}_{\Gamma}(\mathbf{j} \cdot \mathbf{m}). Clearly \mathbf{p} = \frac{1}{\text{vol}(\mathbf{j} \cdot \mathbf{m})} \mathbf{c}(\alpha(\mathbf{j} \cdot \mathbf{m})). Choose \mathbf{j}_q \to \infty to deduce \mathbf{p} \in \Pi_{\Gamma}. Let \beta \in \widetilde{C}_{\Gamma}(\langle \mathbf{n} \rangle). So \mathbf{q} := \frac{1}{\text{vol}(\mathbf{n})} \mathbf{c}(\beta) \in \Pi_{\Gamma}.
```

Claim: For
$$i, j \in \mathbb{N}$$
 $\frac{i}{i+i}\mathbf{p} + \frac{j}{i+i}\mathbf{q} \in \Pi_{\Gamma}$.

Let $\alpha(\mathbf{n} \cdot \mathbf{m}), \beta(\mathbf{m} \cdot \mathbf{n}) \in C_{\Gamma}(\mathbf{n} \cdot \mathbf{m})$ defined as above. Let $\mathbf{j} := (m_1 n_1, \dots, m_{d-1} n_{d-1}, (i+j) m_d n_d)$ view box $\langle \mathbf{j} \rangle$ composed of i+j boxes isomorphic to $\langle \mathbf{m} \cdot \mathbf{n} \rangle$ aligned side-by-side along the direction of \mathbf{e}_d . Color the first i of these boxes by $\alpha(\mathbf{m} \cdot \mathbf{n})$ and the last j by $\beta(\mathbf{n} \cdot \mathbf{m})$, to get $\gamma \in \widetilde{C}_{\Gamma}(\langle \mathbf{j} \rangle)$ with $\frac{1}{\text{vol}(\mathbf{j})}\mathbf{c}(\gamma) = \frac{i}{i+j}\mathbf{p} + \frac{j}{i+j}\mathbf{q}$. Hence $\frac{i}{i+j}\mathbf{p} + \frac{j}{i+j}\mathbf{q} \in \Pi_{\Gamma}$. Since Π_{Γ} is closed $a\mathbf{p} + (1-a)\mathbf{q} \in \Pi_{\Gamma}$ for all $a \in [0,1]$.

Outline of proof-II

Let $\widetilde{\Pi}_{\Gamma}$ be the convex hull of $\frac{1}{\text{vol(m)}}\mathbf{c}(\alpha)$ for some \mathbf{m} and some $\alpha \in \widetilde{C}_{\Gamma}(\langle \mathbf{m} \rangle)$. So $\widetilde{\Pi}_{\Gamma} \subseteq \Pi_{\Gamma}$. The padding part of definition of Γ friendly implies $\widetilde{\Pi}_{\Gamma} \subseteq \Pi_{\Gamma} \subseteq \operatorname{Cl} \widetilde{\Pi}_{\Gamma} \Rightarrow \Pi_{\Gamma} = \operatorname{Cl} \widetilde{\Pi}_{\Gamma}$ Equality $\Pi_{\Gamma} = \operatorname{dom} P_{\Gamma}^*$ follows from last part of Thm 1.

Outline of proof-III

(b) The padding part of definition of Γ friendly implies

For
$$\mathbf{p}, \mathbf{q} \in \Pi_{\Gamma}, \varepsilon > 0 \exists$$

$$\mathbf{m}_{q} := (m_{1,q}, \dots, m_{d,q}), \mathbf{n}_{q} := (n_{1,q}, \dots, n_{d,q}) \in \mathbb{N}^{d}, q \in \mathbb{N}, \ \mathbf{m}_{q}, \mathbf{n}_{q} \to \infty$$
s.t.

$$\begin{split} &\widetilde{C}_{\Gamma}(\langle \mathbf{m}_q \rangle, \mathbf{c}_q), \widetilde{C}_{\Gamma}(\langle \mathbf{n}_q \rangle, \mathbf{d}_q) \neq \emptyset, q \in \mathbb{N}, \\ &\lim_{\mathbf{m}_q \to \infty} \frac{1}{\text{vol}(\mathbf{m}_q)} \mathbf{c}_q = \mathbf{p}, \lim_{\mathbf{n}_q \to \infty} \frac{1}{\text{vol}(\mathbf{n}_q)} \mathbf{d}_q = \mathbf{q}, \\ &\lim_{q \to \infty} \frac{\log \# \widetilde{C}_{\Gamma}(\langle \mathbf{m}_q \rangle, \mathbf{c}_q)}{\text{vol}(\mathbf{m}_q)} \geq h_{\Gamma}(\mathbf{p}) - \varepsilon, \\ &\lim_{q \to \infty} \frac{\log \# \widetilde{C}_{\Gamma}(\langle \mathbf{n}_q \rangle, \mathbf{d}_q)}{\text{vol}(\mathbf{n}_q)} \geq h_{\Gamma}(\mathbf{q}) - \varepsilon. \end{split}$$

Outline of proof - IV

```
Observation that for any \mathbf{m}, \mathbf{n} \in \mathbb{N}^d, \mathbf{c} \in \Pi_n(\text{vol}(\mathbf{m})):
\#\widetilde{C}_{\Gamma}(\langle \mathbf{n} \cdot \mathbf{m} \rangle, \text{vol}(\mathbf{n})\mathbf{c}) > (\#\widetilde{C}_{\Gamma}(\langle \mathbf{m} \rangle, \mathbf{c}))^{\text{vol}(\mathbf{n})} \text{ vields}:
For i, j \in \mathbb{N} h_{\Gamma}(\frac{i}{i+i}\mathbf{p} + \frac{j}{i+i}\mathbf{q}) \geq \frac{i}{i+i}h_{\Gamma}(\mathbf{p}) + \frac{j}{i+i}h_{\Gamma}(\mathbf{q}) - \varepsilon
which proves the concavity of h_{\Gamma}.
(c-d): Let \mathbf{u} \in \operatorname{diff} P_{\Gamma}.
Then \Pi_{\Gamma}(\mathbf{u}) = {\nabla P_{\Gamma}(\mathbf{u})} = \partial P_{\Gamma}(\mathbf{u}) and (c-d) trivially hold.
Recall S(\mathbf{u}) \subseteq \Pi_{\Gamma}(\mathbf{u}), conv S(\mathbf{u}) = \partial P_{\Gamma}(\mathbf{u}) \supset \Pi_{\Gamma}(\mathbf{u}) Let
\mathbf{p}_i \in S(\mathbf{u}), i = 1, \dots, j. So P_{\Gamma}(\mathbf{u}) = \mathbf{p}_i^{\top} \mathbf{u} + h_{\Gamma}(\mathbf{p}_i), i = 1, \dots, j
Since \Pi_{\Gamma} convex, for \mathbf{a} = (a_1, \dots, a_i) \in \Pi_i \mathbf{p} := \sum_{i=1}^J a_i \mathbf{p}_i \in \Pi_{\Gamma}.
As h_{\Gamma} concave P_{\Gamma}(\mathbf{u}) = \sum_{i=1}^{J} a_i \mathbf{p}_i^{\top} \mathbf{u} + h_{\Gamma}(\mathbf{p}_i) < \mathbf{p}^{\top} \mathbf{u} + h_{\Gamma}(\mathbf{p})
The maximal characterization of P_{\Gamma}(\mathbf{u}) implies P_{\Gamma}(\mathbf{u}) = \mathbf{p}^{\top}\mathbf{u} + h_{\Gamma}(\mathbf{p}).
So \mathbf{p} \in \Pi_{\Gamma}(\mathbf{u}) and h_{\Gamma}(\mathbf{p}) = \sum_{i=1}^{J} a_i h_{\Gamma}(\mathbf{p}_i).
(e) Follows from Thm 1 and extra arguments using convexity of P_{\Gamma}^*
```

Reduction of one parameter

```
P_{\Gamma}(\mathbf{u}) = t + P_{\Gamma}(\mathbf{u} - t\mathbf{1}) \Rightarrow \partial P_{\Gamma}(\mathbf{u}) \in \Pi_n It is enough to compute \hat{P}_{\Gamma}(\hat{\mathbf{u}}) := P_{\Gamma}(\hat{\mathbf{u}}), \hat{\mathbf{u}} = (u_1, \dots, u_{n-1}, 0) Hard core model: \hat{P}_{\Gamma}(t) depends on the energy t \in \mathbb{R}. (It is known that for d \geq 2 hard core model has phase transition) For the dimer problem the pressure P_d(\mathbf{v}) depends on \mathbf{v} = (v_1, \dots, v_d), where v_i is the energy of the dimer in the direction \mathbf{e}_i, i = 1, \dots (Non-isotropic model) Dimer isotropic model in \mathbb{Z}^d:
```

pressure $P_d(v)$, where v is the energy of the dimer in any direction.

(Standard model-No phase transition for $v \in \mathbb{R}$)

Computation of pressure

Using the scaled transfer matrices on the torus

 $T(\mathbf{m}'), \mathbf{m}' = (m_1, \dots, m_{d-1})$ as in Friedland-Peled 2005 [3].

Assume for simplicity d = 2, $\Gamma = (\Gamma_1, \Gamma_2)$, where Γ_1 symmetric digraph.

Let Δ transfer digraph induced by Γ_2 between the allowable Γ_1 coloring of the circle T(m).

Then $V := C_{\Gamma_1, \text{per}}(m)$ are the set of vertices of $\Delta(m)$. For

 $\alpha, \beta \in C_{\Gamma_1, \text{per}}(m)$ the directed edge (α, β) is in $\Delta(m)$ iff the

configuration $[(\alpha, \beta)]$ is an allowable configuration on $C_{\Gamma}((m, 2))$.

Adjacency matrix $D(\Delta(m)) = (d_{\alpha\beta})_{\alpha,\beta \in C_{\Gamma_1,per}(m)}$ is $N \times N$ matrix, where

 $N := \#C_{\Gamma_1, \text{per}}(m).$

One dimensional SOFT is $C_{\Gamma}(T(m) \times \mathbb{Z})$:

all Γ allowable coloring of the infinite torus in the direction \mathbf{e}_2 with the basis T(m).

The pressure corresponding to this one dimensional SOFT is denoted by $\tilde{P}_{\Delta(m)}(\mathbf{u})$. Its formula:



Computation of pressure II

Let
$$\tilde{D}(\Delta(m), \mathbf{u}) = (\tilde{d}_{\alpha\beta}(\mathbf{u}))_{\alpha,\beta \in C_{\Gamma_1,\mathrm{per}}(m)} \tilde{d}_{\alpha\beta}(\mathbf{u}) = d_{\alpha\beta} e^{\frac{1}{2}(\mathbf{c}(\alpha) + \mathbf{c}(\beta))^{\top} \mathbf{u}}$$

Then $\tilde{P}_{\Delta}(\mathbf{u}) := \frac{\theta(\mathbf{u},m)}{m}, \theta(\mathbf{u},m) := \log \rho(\tilde{D}(\Delta(m),\mathbf{u}))$
We divide $\log \rho(\tilde{D}(\Delta,\mathbf{u}))$ by m , to have $\tilde{P}_{\Delta}(\mathbf{u} + t\mathbf{1}) = \tilde{P}_{\Delta}(\mathbf{u}) + t$ for any $t \in \mathbb{R}$
Main inequalities
$$\frac{1}{\rho}(\theta(\mathbf{u}, \rho + 2q) - \theta(\mathbf{u}, 2q)) \leq P_{\Gamma}(\mathbf{u}) \leq \frac{1}{2m}(\theta(\mathbf{u}, 2m))$$
for any $m, p > 1$ and $q > 0$.

Automorphism Subgroups

$$A = (a_{ij})_1^N \text{ nonnegative matrix}$$

$$A(A) := \{ \pi \in S_N : a_{\pi(i)\pi(j)} = a_{ij}, i, j \in < N > \}$$
Let $G \leq A(A), \mathcal{O}(G) := < N > /G, M = \#\mathcal{O}(G)$

$$\hat{A} = (\hat{a}_{\alpha\beta})_{\alpha,\beta\in\mathcal{O}(G)}, \hat{a}_{\alpha\beta} =: \sum_{j\in\beta} a_{ij}, i \in \alpha, \rho(A) = \rho(\hat{A}),$$
If $A = A^T$ then \hat{A} symmetric for
$$< \mathbf{x}, \mathbf{y} >= \sum_{\alpha\in\mathcal{O}(G)} (\#\alpha) \mathbf{x}_{\alpha} \mathbf{y}_{\alpha}. M \geq N/\#G,$$
In our computations $M \sim N/\#G$

Using these tools we confirmed Baxter's computations with nine digits of precision of $P_2(v)$ and of $h_2(p)$.

We also computed the non-isotropic $P_2((v_1, v_2))$.

Graph 1 of pressure 2-dimensional dimers $P(v_1, v_2)$

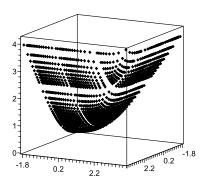


Figure: The graph of $\frac{\bar{P}_1(12,(v_1,v_2))}{12}$ for angle $\theta=28^{\rm o}, \varphi=78^{\rm o}$

Graph 2 of pressure 2-dimensional dimers $P(v_1, v_2)$

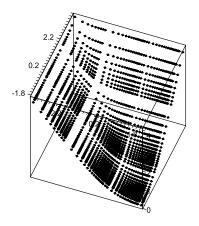


Figure: The graph of $\frac{\bar{P}_1(12,(v_1,v_2))}{12}$ for angle $\theta=-159^{\circ}, \varphi=42^{\circ}$

Graph 1 of 2-dimensional entropy $h_2(v_1, v_2)$

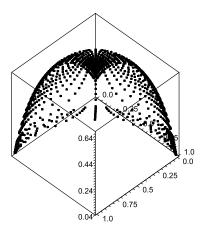


Figure: The graph of an approximation of $\bar{h}_2(p_1, p_2)$ for angle $\theta = 45^{\circ}, \varphi = 45^{\circ}$

Graph 2 of 2-dimensional entropy $h_2(v_1, v_2)$

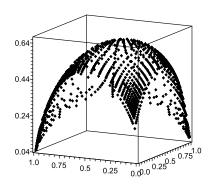


Figure: $\bar{h}_2(p_1, p_2)$ for angle $\theta = -153^{\circ}, \varphi = 78^{\circ}$

- V. I. Arnold, *Mathematical Methods of Classical Mechanics*, Spinger, New York, 1978.
- H. Au-Yang and J.H.H. Perk, Phys. Lett. A 104 (1984), 131–134.
- R. J. Baxter, Dimers on a rectangular lattice, *J. Math. Phys.* 9 (1968), 650–654.
- S. Friedland, On the entropy of Z-d subshifts of finite type, *Linear Algebra Appl.* 252 (1997), 199–220.
- S. Friedland, Multi-dimensional capacity, pressure and Hausdorff dimension, in *Mathematical System Theory in Biology, Communication, Computation and Finance*, edited by J. Rosenthal and D. Gilliam, IMA Vol. Ser. 134, Springer, New York, 2003, 183–222.
- S. Friedland and L. Gurvits, Lower bounds for partial matchings in regular bipartite graphs and applications to the monomer-dimer entropy, Combinatorics, Probability and Computing, 17 (2008), 347-361.

- S. Friedland, E. Krop, P. H. Lundow and K. Markström, On the validations of the asymptotic matching conjectures, *Journal of Statistical Physics*, 133 (2008), 513-533.
- S. Friedland, E. Krop and K. Markstrom, On the number of matchings in regular graphs, *The Electronic Journal of Combinatorics*, 15 (2008), #R110, 1-28.
- S. Friedland and U. N. Peled, Theory of Computation of Multidimensional Entropy with an Application to the Monomer-Dimer Problem, *Advances of Applied Math.* 34(2005), 486-522.
- F. R. Gantmacher, *The Theory of Matrices*, Vol. II, Chelsea Publ. Co., New York 1959.
- J. M. Hammersley, Existence theorems and Monte Carlo methods for the monomer-dimer problem, in *Reseach papers in statistics:* Festschrift for J. Neyman, edited by F.N. David, Wiley, London, 1966, 125–146.

- O. J. Heilman and E. H. Lieb, Theory of monomer-dimer systems, *Comm. Math. Phys.* 25 (1972), 190–232; Errata 27 (1972), 166.
- E. Ising, Beitrag zur Theory des Ferromagnetismus, *Z. Physik* 31 (1925), 253–258.
- P.W. Kasteleyn, The statistics of dimers on a lattice, *Physica* 27 (1961), 1209–1225.
- J. F. C. Kingman, A convexity property of positive matrices. *Quart. J. Math. Oxford Ser.* (2) 12 (1961), 283–284.
- L. Onsager, Cristal statistics, I. A two-dimensional model with an order-disorder transition, *Phys. Review* 65 (1944), 117–149.
- R. Peierles, On Ising model of ferrogmanetism, *Proc. Cambridge Phil. Soc.* 32 (1936), 477-481.
- R.B. Potts, ?, Proc. Cambridge Phil. Soc. 48 (1952), 106-?
- R. T. Rockafeller, Convex Analysis, Princeton Univ. Press 1970.

- H. N. V. Temperley and M. E. Fisher, Dimer probelm in statistical mechanics—an exact result, *Phil. Mag. (8)* 6 (1961), 1061–1063.
- C. J. Thompson, *Mathematical Statisitical Mechanics*, Princeton Univ. Press, 1972.